# Isobarlc Heat Capacity Data for Liquid HCFC-123 $\left(\mathrm{CHCl}_{2} \mathrm{CF}_{3}\right.$, 2,2-Dichloro-1,1,1-trifluoroethane) 

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The lsobaric heat capactiy, $C_{p}$, of Ilquid
2,2-dichloro-1,1,1-tritituoroethane ( $\mathrm{CHCl}_{2} \mathrm{CF}_{3}, \mathrm{HCFC}-123$ ), which is a promising alternative to CFC-11, has been measured by using fiow calorinetry. The values of $80 C_{p}$ have been determined in the range of temperatures from 275.65 to $\mathbf{4 4 0} \mathrm{K}$ and pressures from 0.5 to 3.2 MPa , respectively. The overall uncertainties of the defermined $C_{p}$ values are concluded to be less than $\pm 0.4 \%$ for temperatures below $420 \mathrm{~K}, \pm 0.5 \%$ at $\mathbf{4 3 0} \mathrm{K}, \pm 0.8 \%$ at 440 K , respectively. The $C_{p}$ data have been correlated whit a function of temperature and pressure within $\pm 0.4 \%$, and the $C_{p}$ of saturated IIquid, $C_{p}{ }^{\prime}$, have been derived from the correlation.

## Introduction

The fully halogenated chiorofluorocarbons (CFCs) have been widely used as a blowing agent, a cleaning agent, or a working fluld for heat-pumping and refrigeration systems. But there is concern for their ozone-depletion and global-warming potentlal so that many CFC alternatives have been suggested. HCFC$123\left(\mathrm{CHCl}_{2} \mathrm{CF}_{3}\right)$, which has small ozone-depletion potentlal, is promised as one of alternatives to replace CFC-11 ( $\mathrm{CCl}_{3} \mathrm{~F}$ ). This paper reports the isobaric heat capacity data of liquid HCFC-123 and a correlation for the heat capacity of compressed and saturated liquid. The purity of the sample HCFC-123 used in the measurements was 99.82 wt $\%$.

## Experlmental Section

The detailed descriptlon regarding the flow calorimeter has been reported in our previous papers (1-3). Measurements for another alternative refrigerant, HFC -134a $\left(\mathrm{CF}_{3} \mathrm{CH}_{2} \mathrm{~F}\right)$, have already been reported by the present authors (4). The speclal features of our apparatus are lis highly adlabatic performance, the flow stability of the sample liquid in the closed circulation system, and the rellability of automatic measurements of the mass-flow rate. The isobaric heat capacity, $C_{p}$, is defined as follows:

$$
\begin{equation*}
C_{p}=\dot{Q} /(\dot{m} \Delta T) \tag{1}
\end{equation*}
$$

Flow calorimetry consists of three simultaneous measurements: measurement of energy, $\dot{Q}$, supplied by a microheater to the flowing sample liquid; measurement of the temperature increment, $\Delta T$, which is the temperature difference, $T_{\text {out }}-T_{m}$, of sample liquid before and after heating by the microheater; and measurement of the mass-flow rate, $\dot{m}$. The reliability of this apparatus was confirmed by measuring the $C_{p}$ values of water. The systematic errors were not found, and the standard deviation from the equation of state developed by Sato et al. (5) was $0.34 \%$.

## Results

Measurements were performed at temperatures from 275.65 to 440 K and pressures from 0.5 to 3.2 MPa . All measured values are listed in Table I. The table includes the measured

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pressure, $P$, measured temperature, $T$, energy supplied to the sample fluid, $\dot{Q}$, temperature increment, $\Delta T$, mass-flow rate, $\dot{m}$, and measured heat capacity, $C_{p}$. Note that the measured temperature $T_{68}$ is assigned to the temperature of the arithmetic mean of $T_{\text {in }}$ and $T_{\text {out }}$ on the basis of the International Practical Temperature Scale of 1968 (IPTS-68).

In order to know the effect of heat loss on measured $C_{p}$ values, data at varlous mass-flow rates are plotted against the inverse mass-flow rates in Figure 1. Measurements were performed at a state, i.e., at a certain temperature and pressure, selecting two different mass-flow rates. It was confirmed from Figure 1 that the measured $C_{p}$ values do not depend on the mass-flow rates. Thus, it was concluded that the effect of heat loss is small enough to not necessarily be compensated for. The temperature increment has to be small enough so as to produce no difference between the average of the $C_{p}$ values at inlet and outlet temperatures, $C_{p}(a v)$, and the true $C_{p}$ value at a given temperature. The measurements in the temperature range from 275.65 to 420 K were performed with the temperature increment 5 K , while the measurements at 430 K were performed with about 3 K of $\Delta T$ and 1.6 K for the measurements at 440 K , respectively. The unsmoothed experimental data are summarized in Table I. We confirmed that the difference between $C_{p}$ (av) and the true $C_{p}$ was within $\pm 0.1 \%$ by examining the data with a developed correlation given as a function of temperature and pressure. After we corrected the $C_{p}$ values that were measured at the same temperatures and pressures but mass-flow rates to those at nominal temperatures and pressures with the ald of developed correlation, we determined the $C_{p}$ values at nominal temperatures and pressures as arithmetic means of those at the same temperatures and pressures but mass-flow rate. In Table II, $80 C_{p}$ values at nominal temperatures and pressures are listed. The uncertainties of the measurements are $\pm 8 \mathrm{mK}$ in temperature increment, $\pm 0.01 \%$ in energy supplied, $\pm 0.3 \%$ in mass-flow rate, $\pm 11 \mathrm{mK}$ in temperature, and $\pm 3 \mathrm{kPa}$ in pressure, respectively. The overall uncertainties of the determined $C_{p}$ values summarized in Table II is concluded to be less than $\pm 0.4 \%$ for temperatures below $420 \mathrm{~K}, \pm 0.5 \%$ for those at 430 K , and $\pm 0.8 \%$ for those at 440 K , respectively. The $C_{p}$ values were correlated with the following temperature and pressure function.

$$
\begin{gather*}
C_{p} / R=a+b P_{r}^{0.5}+c P_{r}  \tag{2}\\
a=0.06718\left(1-T_{r}\right)^{-2}-0.3756\left(1-T_{r}\right)^{-1}+ \\
27.24-12.98\left(1-T_{r}\right)^{0.5} \\
b=6.215 \times 10^{-7}\left(1-T_{r}\right)^{-5}-0.001719\left(1-T_{r}\right)^{-3}- \\
0.6750\left(1-T_{r}\right) \\
c=-0.01581\left(1-T_{r}\right)^{-2}+0.5316\left(1-T_{r}\right)
\end{gather*}
$$

where $P_{r}=P / P_{c}, T_{r}=T / T_{c}, R=R_{0} / M$, and $C_{p}$ is given in $\mathrm{kJ} /(\mathrm{kg} \cdot \mathrm{K}), P$ in MPa, and $T$ in K. The critical pressure, $P_{\mathrm{c}}$, is 3.6655 MPa , which has been reported by Plao et al. (6). The critical temperature, $T_{c}$, is 456.86 K , which has been reported by Tanikawa et al. (7). The universal gas constant $R_{0}=$ $8.31451 \mathrm{~J} /(\mathrm{mol} \cdot \mathrm{K})$ and molar mass $M=152.93 \mathrm{~g} / \mathrm{mol}$.

Table I. Measured Ieobaric Heat Capacity of HCFC-123

| $P, \mathrm{MPa}$ | T, K | Q, J/s | $\Delta T, \mathrm{~K}$ | m, g/s | $\begin{gathered} C_{p}, \\ \mathrm{~kJ} /(\mathrm{kg} \cdot \mathrm{~K}) \end{gathered}$ | $P, \mathrm{MPa}$ | T, K | $\dot{Q}, \mathrm{~J} / \mathrm{s}$ | $\Delta T, \mathrm{~K}$ | m, g/s | $\begin{gathered} C_{p,} \\ \mathrm{~kJ} /(\mathrm{kg} \cdot \mathrm{~K}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.500 | 275.64 | 0.8177 | 4.998 | 0.1623 | 1.008 | 0.802 | 360.00 | 1.0182 | 5.020 | 0.1791 | 1.133 |
| 0.503 | 275.65 | 0.5649 | 5.009 | 0.1123 | 1.004 | 0.802 | 360.00 | 0.8290 | 5.011 | 0.1462 | 1.132 |
| 1.004 | 275.65 | 0.8150 | 5.000 | 0.1627 | 1.002 | 1.002 | 360.00 | 1.0179 | 5.021 | 0.1794 | 1.130 |
| 1.004 | 275.65 | 0.5629 | 5.013 | 0.1123 | 1.000 | 1.002 | 359.99 | 0.8266 | 5.005 | 0.1461 | 1.131 |
| 1.501 | 275.65 | 0.8167 | 5.005 | 0.1631 | 1.000 | 1.502 | 359.99 | 1.0142 | 5.026 | 0.1792 | 1.126 |
| 1.501 | 275.65 | 0.5629 | 5.004 | 0.1127 | 0.998 | 1.503 | 360.00 | 0.8263 | 5.008 | 0.1464 | 1.127 |
| 2.002 | 275.66 | 0.8168 | 5.016 | 0.1625 | 1.002 | 2.000 | 360.00 | 1.0127 | 5.016 | 0.1799 | 1.122 |
| 2.001 | 275.66 | 0.5625 | 5.019 | 0.1123 | 0.998 | 2.000 | 360.00 | 0.8241 | 5.017 | 0.1461 | 1.124 |
| 2.499 | 275.65 | 0.8160 | 5.009 | 0.1631 | 0.999 | 2.499 | 359.99 | 1.0059 | 5.002 | 0.1798 | 1.119 |
| 2.500 | 275.66 | 0.5648 | 5.018 | 0.1129 | 0.997 | 2.499 | 360.00 | 0.8230 | 5.017 | 0.1462 | 1.122 |
| 3.000 | 275.65 | 0.8172 | 5.006 | 0.1631 | 1.001 | 2.998 | 360.00 | 0.9117 | 5.013 | 0.1634 | 1.113 |
| 2.998 | 275.68 | 0.5709 | 5.070 | 0.1128 | 0.998 | 1.003 | 370.00 | 1.0376 | 5.021 | 0.1799 | 1.149 |
| 0.501 | 300.01 | 0.8453 | 5.020 | 0.1635 | 1.030 | 1.003 | 370.00 | 0.8431 | 5.017 | 0.1463 | 1.149 |
| 0.503 | 300.00 | 0.5835 | 5.011 | 0.1130 | 1.030 | 1.502 | 370.00 | 1.0336 | 5.017 | 0.1799 | 1.145 |
| 1.002 | 300.00 | 0.8401 | 5.010 | 0.1631 | 1.029 | 1.502 | 370.00 | 0.8403 | 5.017 | 0.1465 | 1.143 |
| 1.000 | 300.00 | 0.5803 | 5.009 | 0.1126 | 1.029 | 2.005 | 370.02 | 1.0303 | 5.036 | 0.1798 | 1.138 |
| 1.502 | 300.00 | 0.8389 | 5.003 | 0.1632 | 1.027 | 2.010 | 370.01 | 0.8308 | 5.013 | 0.1457 | 1.137 |
| 1.500 | 299.99 | 0.5787 | 4.998 | 0.1129 | 1.026 | 2.500 | 370.00 | 1.0305 | 5.034 | 0.1805 | 1.134 |
| 2.000 | 300.00 | 0.8399 | 5.006 | 0.1632 | 1.028 | 2.498 | 370.00 | 0.8347 | 5.017 | 0.1467 | 1.134 |
| 1.999 | 300.00 | 0.5806 | 5.009 | 0.1131 | 1.025 | 3.000 | 369.99 | 1.0369 | 5.012 | 0.1829 | 1.131 |
| 2.505 | 300.00 | 0.8394 | 5.008 | 0.1636 | 1.025 | 2.998 | 369.99 | 0.8440 | 5.003 | 0.1490 | 1.132 |
| 2.506 | 299.99 | 0.5812 | 5.000 | 0.1136 | 1.023 | 1.200 | 380.01 | 1.0537 | 5.040 | 0.1777 | 1.176 |
| 3.001 | 300.00 | 0.8432 | 5.013 | 0.1641 | 1.025 | 1.199 | 380.01 379.97 | 0.8588 | 4.040 | 0.1463 | 1.178 |
| 2.999 | 300.00 | 0.5826 | 5.027 | 0.1133 | 1.023 | 1.499 | 379.99 | 1.0538 | 4.983 5.025 | 0.1789 | 1.172 |
| 0.504 | 310.00 | 0.8498 | 5.004 | 0.1625 | 1.045 | 1.500 | 380.00 | 0.8604 | 5.015 | 0.1462 | 1.173 |
| 0.504 | 310.00 | 0.5867 | 4.994 | 0.1123 | 1.047 | 2.000 | 380.01 | 0.8603 | 5.020 | 0.1473 | 1.164 |
| 1.002 | 310.01 | 0.8528 | 5.027 | 0.1627 | 1.043 | 2.003 | 380.02 | 1.0530 | 5.040 | 0.1794 | 1.165 |
| 1.003 | 310.00 | 0.5898 | 5.006 | 0.1129 | 1.044 | 2.502 | 380.02 | 1.0535 | 5.057 | 0.1800 | 1.157 |
| 1.505 | 310.00 | 0.8489 | 5.008 | 0.1624 | 1.044 | 2.500 | 380.00 | 0.8546 | 5.035 | 0.1469 | 1.155 |
| 1.505 | 310.00 | 0.5882 | 4.998 | 0.1125 | 1.046 | 2.997 | 380.00 | 1.0556 | 5.030 | 0.1821 | 1.152 |
| 2.005 | 309.99 | 0.8475 | 4.992 | 0.1628 | 1.043 | 2.996 | 379.99 | 0.8630 | 5.009 | 0.1494 | 1.154 |
| 2.005 | 310.01 | 0.5876 | 5.005 | 0.1125 | 1.044 | 1.501 | 390.01 | 1.0921 | 5.035 | 0.1797 | 1.207 |
| 2.503 | 310.01 | 0.8485 | 5.013 | 0.1622 | 1.044 | 1.501 | 390.01 | 1.88870 | 5.033 | 0.1458 | 1.208 |
| 2.503 | 310.01 | 0.5854 | 5.006 | 0.1121 | 1.043 | 1.999 | 390.00 | 1.0825 | 5.032 | 0.1796 | 1.198 |
| 3.004 | 310.01 | 0.8506 | 5.009 | 0.1629 | 1.043 | 1.999 | 390.01 | 0.8840 | 5.030 | 0.1470 | 1.196 |
| 3.002 | 310.00 | 0.5879 | 4.994 | 0.1129 | 1.043 | 2.502 | 390.01 | 1.0801 | 5.051 | 0.1801 | 1.187 |
| 0.505 | 320.01 | 0.8637 | 5.020 | 0.1620 | 1.062 | 2.497 | 390.02 | 0.8788 | 5.052 | 0.1463 | 1.189 |
| 0.507 | 320.00 | 0.5970 | 4.996 | 0.1125 | 1.062 | 2.998 | 390.00 | 1.0855 | 5.034 | 0.1828 | 1.180 |
| 1.000 | 320.00 | 0.8622 | 5.000 | 0.1626 | 1.061 | 2.997 | 390.00 | 0.8858 | 5.033 | 0.1489 | 1.182 |
| 1.002 | 320.00 | 0.5971 | 5.009 | 0.1122 | 1.062 | 1.702 | 400.00 | 1.1314 | 5.045 | 0.1804 | 1.243 |
| 1.504 | 320.00 | 0.8573 | 4.998 | 0.1618 | 1.060 | 1.703 | 400.01 | 0.9217 | 5.041 | 0.1468 | 1.245 |
| 1.504 | 320.00 | 0.5927 | 4.996 | 0.1118 | 1.061 | 1.997 | 400.00 | 1.1203 | 5.043 | 0.1800 | 1.234 |
| 2.001 | 320.00 | 0.8598 | 5.003 | 0.1628 | 1.056 | 1.998 | 399.99 | 0.9101 | 5.010 | 0.1467 | 1.238 |
| 2.001 | 320.00 320.01 | 0.5958 0.8623 | 4.997 5.013 | 0.1128 0.1630 | 1.057 1.055 | 2.501 | 400.00 | 1.1119 | 5.045 | 0.1802 | 1.223 |
| 2.500 | 320.01 | 0.8623 0.5975 | 5.013 5.018 | 0.1630 0.1127 | 1.055 1.056 | 2.498 | 400.00 | 0.9065 | 5.040 | 0.1469 | 1.224 |
| 2.998 | 320.00 | 0.8719 | 4.996 | 0.1656 | 1.054 | 3.002 | 399.98 | 0.9095 | 4.991 | 0.1497 | 1.217 |
| 3.001 | 320.00 | 0.6086 | 5.003 | 0.1156 | 1.052 | 3.004 | 399.99 | 1.0975 | 5.034 | 0.1793 | 1.216 |
| 0.502 | 330.00 | 0.8731 | 5.000 | 0.1623 | 1.076 | 1.903 | 410.02 | 1.1799 | 5.074 | 0.1788 | 1.300 |
| 0.503 | 330.00 | 0.6064 | 5.002 | 0.1123 | 1.080 | 1.903 | 410.01 | 0.9648 | 5.054 | 0.1462 | 1.305 |
| 1.003 | 330.00 | 0.8717 | 5.002 | 0.1622 | 1.074 | 1.998 1.998 | 410.01 410.02 | 1.1765 1.1765 | 5.055 5.063 | 0.1796 | 1.296 |
| 1.001 | 330.01 | 0.6065 | 5.019 | 0.1122 | 1.077 | 1.998 | 410.02 410.01 | 1.1765 1.1765 | 5.063 5.039 | 0.1803 | 1.295 |
| 1.505 | 330.00 | 0.8713 | 5.004 | 0.1628 | 1.070 | 1.998 | 410.01 | 0.9609 | 5.039 | 0.1473 | 1.294 |
| 1.503 | 330.00 | 0.6035 | 4.996 | 0.1125 | 1.073 | 1.998 | 410.01 | 0.9609 | 5.039 | 0.1471 | 1.296 |
| 2.002 | 330.00 330.00 | 0.8702 0.6030 | 5.008 | 0.1622 | 1.072 | 1.998 | 410.01 | 0.9609 | 5.041 | 0.1470 | 1.297 |
| 2.004 | 330.00 | 0.6030 | 4.999 | 0.1123 | 1.074 | 2.501 | 410.00 | 1.1624 | 5.064 | 0.1802 | 1.274 |
| 2.496 | 330.01 | 0.8785 | 5.019 5.039 | 0.1631 | 1.073 | 2.499 | 410.01 | 0.9468 | 5.057 | 0.1465 | 1.278 |
| 2.496 2.999 | 330.02 330.00 | 0.6108 0.8707 | 5.039 5.012 | 0.1130 0.1626 | 1.073 | 3.000 | 410.00 | 1.1560 | 5.041 | 0.1829 | 1.254 |
| 2.999 | 330.00 | 0.8707 0.6173 | 5.012 | 0.1626 | 1.068 | 3.000 | 410.00 | 1.1560 | 5.043 | 0.1829 | 1.253 |
| 3.003 | 330.00 | 0.6173 | 4.997 | 0.1154 | 1.070 | 3.000 | 410.01 | 0.9519 | 5.060 | 0.1496 | 1.257 |
| 0.603 | 339.99 | 0.8845 | 5.001 | 0.1621 | 1.091 | 2.302 | 420.02 | 1.2465 | 5.070 | 0.1791 | 1.373 |
| 1.004 | 340.00 | 0.8862 | 5.012 | 0.1624 | 1.089 | 2.499 | 420.00 | 1.2465 1.2308 | 5.044 | 0.1795 | 1.359 |
| 1.501 | 340.01 | 0.8872 | 5.014 | 0.1626 | 1.088 | 2.994 | 420.01 | 1.2071 | 5.048 | 0.1801 | 1.328 |
| 2.001 | 340.00 | 0.8861 | 5.008 | 0.1633 | 1.083 | 3.194 | 420.00 | 1.1950 | 5.045 | 0.1800 | 1.316 |
| 2.503 | 340.00 | 0.8857 | 5.010 | 0.1632 | 1.083 | 3.200 | 420.01 | 1.3089 | 5.069 | 0.1960 | 1.317 |
| 3.001 | 340.00 | 0.8991 | 5.010 | 0.1660 | 1.081 |  |  |  |  |  |  |
| 0.702 | 350.00 | 0.9048 | 5.011 | 0.1631 | 1.107 | 2.701 2.803 | 430.00 429.98 | 0.8112 0.8031 | 3.055 3.050 | 0.1799 0.1797 | 1.476 1.465 |
| 1.002 | 349.99 | 0.9049 | 5.005 | 0.1628 | 1.111 | 2.899 | 429.98 429.99 | 0.8031 0.7953 | 3.055 | 0.1802 | 1.445 |
| 1.502 | 350.00 | 0.9045 | 5.019 | 0.1633 | 1.104 | 3.193 | 429.99 | 0.7795 | 3.045 | 0.1802 | 1.421 |
| 2.003 | 349.99 | 0.9004 | 4.993 | 0.1633 | 1.104 | 3.193 3.196 | 429.99 | 0.8469 | 3.045 3.046 | 0.1956 | 1.421 |
| 2.499 | 350.00 | 0.9000 | 5.018 | 0.1634 | 1.098 | 3.196 | 429.99 | 0.8469 | 3.046 | 0.1956 | 1.421 |
| 3.001 | 350.00 | 0.9127 | 5.001 | 0.1662 | 1.098 | 3.000 | 440.02 | 0.4887 | 1.582 | 0.1793 | 1.723 |
|  |  |  |  |  |  | 3.199 | 440.01 | 0.4665 | 1.577 | 0.1800 | 1.644 |
|  |  |  |  |  |  | 3.200 | 440.01 | 0.5043 | 1.570 | 0.1953 | 1.644 |



Figure 1. Reproducibility in $C_{p}$ measurements at various mass-flow rates. $C_{p}(\exp )$ denotes the measured value, whereas $C_{p}(a v)$ is the averaged value.

Table II. Isobaric Heat Capacity of HCFC-123

| $P, \mathrm{MPa}$ | T, K | $\frac{C_{\mathrm{p}}^{\prime}}{\mathrm{kJ} / \mathrm{kg} \cdot \mathrm{~K})}$ | $P, \mathrm{MPa}$ | T, K | $\underset{\mathrm{kJ} /(\mathrm{kg} \cdot \mathrm{~K})}{C_{\mathrm{p}}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.500 | 275.65 | 1.006 | 0.800 | 360.00 | 1.132 |
| 1.000 | 275.65 | 1.001 | 1.000 | 360.00 | 1.130 |
| 1.500 | 275.65 | 1.000 | 1.500 | 360.00 | 1.127 |
| 2.000 | 275.65 | 1.000 | 2.000 | 360.00 | 1.123 |
| 2.500 | 275.65 | 0.998 | 2.500 | 360.00 | 1.120 |
| 3.000 | 275.65 | 1.000 | 3.000 | 360.00 | 1.113 |
| 0.500 | 300.00 | 1.030 | 1.000 | 370.00 | 1.149 |
| 1.000 | 300.00 | 1.029 | 1.500 | 370.00 | 1.144 |
| 1.500 | 300.00 | 1.027 | 2.000 | 370.00 | 1.138 |
| 2.000 | 300.00 | 1.027 | 2.500 | 370.00 | 1.134 |
| 2.500 | 300.00 | 1.024 | 3.000 | 370.00 | 1.131 |
| 3.000 | 300.00 | 1.024 | 1.200 | 380.00 | 1.177 |
| 0.500 | 310.00 | 1.046 | 1.500 | 380.00 | 1.173 |
| 1.000 | 310.00 | 1.043 | 2.000 | 380.00 | 1.164 |
| 1.500 | 310.00 | 1.045 | 2.500 | 380.00 | 1.156 |
| 2.000 | 310.00 | 1.043 | 3.000 | 380.00 | 1.153 |
| 2.500 | 310.00 | 1.043 |  |  |  |
| 3.000 | 310.00 | 1.043 | 1.500 2.000 | 399.00 390.00 | 1.197 |
| 0.500 | 320.00 | 1.062 | 2.500 | 390.00 | 1.188 |
| 1.000 | 320.00 | 1.061 | 3.000 | 390.00 | 1.181 |
| 1.500 | 320.00 | 1.060 |  |  |  |
| 2.000 | 320.00 | 1.056 | 1.700 2.000 | 400.00 400.00 | 1.244 1.236 |
| 2.500 | 320.00 | 1.056 | 2.500 | 400.00 | 1.223 |
| 3.000 | 320.00 | 1.053 | 2.500 3.000 | 400.00 | 1.223 1.216 |
| $0.500$ | $330.00$ | 1.078 | 1.900 | 410.00 | 1.303 |
| 1.000 | 330.00 | 1.076 | 2.000 | 410.00 | 1.296 |
| 1.500 | 330.00 | 1.071 | 2.500 | 410.00 | 1.276 |
| 2.000 | 330.00 | 1.073 | 3.000 | 410.00 | 1.255 |
| 2.500 | 330.00 | 1.073 |  |  |  |
| 3.000 | 330.00 | 1.069 | 2.300 | 420.00 420 | $\begin{aligned} & 1.373 \\ & 1.350 \end{aligned}$ |
| 0.600 | 340.00 | 1.091 | 2.500 3.000 | 420.00 420.00 | 1.359 1.328 |
| 1.000 | 340.00 | 1.089 | 3.000 3.200 | 420.00 420.00 | 1.328 1.316 |
| 1.500 | 340.00 | 1.088 | 3.200 | 420.00 | 1.316 |
| 2.000 | 340.00 | 1.083 | 2.700 | 430.00 | 1.476 |
| 2.500 | 340.00 | 1.083 | 2.800 | 430.00 | 1.465 |
| 3.000 | 340.00 | 1.081 | 3.000 | 430.00 | 1.445 |
| 0.700 | 350.00 | 1.107 | 3.200 | 430.00 | 1.421 |
| 1.000 | 350.00 | 1.111 | 3.000 | 440.00 | 1.723 |
| 1.500 | 350.00 | 1.104 | 3.200 | 440.00 | 1.644 |
| 2.000 | 350.00 | 1.104 |  |  |  |
| 2.500 | 350.00 | 1.098 |  |  |  |
| 3.000 | 350.00 | 1.098 |  |  |  |

Equation 2 is effective in a temperature range between 275.65 and 440 K . This correlation reproduces the measured $C_{p}$ data within $\pm 0.4 \%$. Saturated liquid $C_{p}$ data in Table III were also derived by substituting vapor pressures calculated from eq 3 ,

$$
\begin{gathered}
\ln \left(P / P_{c}\right)=\left(A \tau+B \tau^{1.2}+C \tau^{2}+D \tau^{3}\right) /(1-\tau) \\
\tau=1-T / T_{c} \\
A=-7.87576 \quad B=1.45751 \\
C=0.520220 \quad D=-3.47970
\end{gathered}
$$

Table III. Isobaric Heat Capacity of Saturated Liquid HCFC-123

|  |  | $C_{\mathrm{p}}{ }^{\prime}{ }^{\prime}$ <br> $\mathrm{kJ} /(\mathrm{kg} \cdot \mathrm{K})$ | $P, \mathrm{MPa}$ | $T, \mathrm{~K}$ | $C_{\mathrm{p}}{ }^{\prime}{ }^{\prime}$ <br> $\mathrm{kJ} /(\mathrm{kg} \cdot \mathrm{K})$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.037 | 275.65 | 1.006 | 0.580 | 360.00 | 1.132 |
| 0.044 | 280.00 | 1.012 | 0.733 | 370.00 | 1.154 |
| 0.067 | 290.00 | 1.024 | 0.914 | 380.00 | 1.180 |
| 0.098 | 300.00 | 1.036 | 1.127 | 390.00 | 1.211 |
| 0.139 | 310.00 | 1.050 | 1.374 | 400.00 | 1.252 |
| 0.193 | 320.00 | 1.064 | 1.659 | 410.00 | 1.308 |
| 0.261 | 330.00 | 1.079 | 1.987 | 420.00 | 1.392 |
| 0.347 | 340.00 | 1.095 | 2.363 | 430.00 | 1.527 |
| 0.452 | 350.00 | 1.112 | 2.792 | 440.00 | 1.800 |



Figure 2. Experimental data and the correlation for isobaric heat capacity of liquid HCFC-123.


Figure 3. Comparison of isobaric heat capacity of saturated liquid HCFC-123 with that of CFC-11 (8).
developed by Plao et al. (6), into eq 2. The measured $C_{p}$ values for the compressed liquid and the isotherms, as well as the saturation curve calculated from eq 2 , are shown on a $C_{p}-P$ diagram in Figure 2.

## Discussion

Since no measured $C_{\rho}$ values nor equations of state that are effective in the liquid phase have been available for HCFC-123 at present, we cannot compare our data with another source. As shown in Figure 1, the reproducibility of our measurements was better than $\pm 0.3 \%$ and our correlation can reproduce our data within $\pm 0.4 \%$. No constrained behavlor beyond the experimental error was observed in our measurements.

For practical applications, when the $C_{\rho}^{\prime}$ values of HCFC-123 are compared with those of CFC-11 (8), the $C_{\rho}$ ' of HCFC-123 is larger than that of CFC-11 by about $10 \%$, as shown in Figure 3.

## Conclusion

The $C_{\rho}$ values of hydrofluorocarbon HCFC-123 were measured at 80 state points in the liquid phase, covering tempera-
tures from 275.65 to 440 K and pressures from 0.5 to 3.2 MPa with the uncertainty of $\pm 0.4 \%$ for the data below 420 K , $\pm 0.5 \%$ for the data at 430 K , and $\pm 0.8 \%$ for the data at 440 $K$, respectively. These $C_{p}$ data were correlated as a function of temperature and pressure, and saturated liquid $C_{p}$ ' values were derlved from the correlation.

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# Vapor-Liquid Equilibrium Determination by Total Pressure Measurements for Three Binary Systems Made of 1,2-Dimethoxyethane with Toluene, Methylcyclohexane, or (Trifluoromethyl)benzene 

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Total vapor pressures were measured for the three blnary systems made of 1,2-dimethoxyethane with toluene, methylcyclohexane, or (trifluoromethyl)benzene (benzotrifluoride) at 350 K . Densities of the mixtures were measured at 298.15 K . The vapor-llquid equillbria were correlated by using the Wilson equation.

## Introduction

Vapor-liquild equilibria (VLE) are required for engineering use such as in the design and operation of distillation equipment. In the present study, isothermal total pressures $P$ were measured by the ebulliometric method for the three binary systems made of 1,2-dimethoxyethane with toluene, methylcyclohexane, or (trifluoromethyl)benzene (benzotrifluoride) at 350 K . The densities of these mixtures were measured at 298.15 K , and the molar excess volumes $V^{E}$ were calculated.

## Experimental Section

The experimental apparatus previously reported by Kato et al. $(1,2)$ was modified in the present study, as shown in Figure 1. The ebulliometer E has been described by Kato et al. (3). The liquid volume in the ebulliometer is about $25 \mathrm{~cm}^{3}$.

At the start of an experiment, a solution of desired composition was prepared by mixing each pure substance, which was weighed by use of syringes and an automatic balance, similarty to the procedures of the prevlous works $(3,4)$. The reproducibility of the composition was within $\pm 0.001$ mole fraction. Cocks $K_{1}, K_{2}$, and $K_{3}$ and solenoid valve $V$ were opened, and $K_{4}$ and $K_{5}$ were closed. The system pressure was reduced by the vacuum pump near to the total pressure of the desired boiling point temperature. Next, cocks $\mathrm{K}_{2}$ and solenoid valve $\checkmark$ were closed, and the temperature was kept constant by use of the Willams (5) two-liquld manostat with a precision of
$\pm 0.01 \mathrm{mmHg}$. The upper layer of the manostat is dibutyl phthalate, and the lower layer is ethylene glycol saturated with sodium nitrite. The prepared solution was then boiled. The temperature was controlled to the desired temperature by reducing the pressure.

After attainment of steady state, the total pressure was measured with a Ruska 3850 quartz Bourdon gage with a precision of $\pm 0.02 \mathrm{mmHg}$. The experimental temperature was measured at 350 K with a Hewlett-Packard 2804A quartz thermometer calibrated at the triple point of water in a reference cell. The reproducibility of the thermometer temperature measurements was $\pm 0.01 \mathrm{~K}$. The densities of the binary liquid mixtures were determined at 298.15 K with an Anton Paar DMA48 digital density meter with a precision of $\pm 0.0001 \mathrm{~g} / \mathrm{cm}^{3}$.

Toluene was a speclal grade reagent supplied by Wako Pure Chemical Industries Ltd. 1,2-Dimethoxyethane, methylcyclohexane, and (trifluoromethyl)benzene were special grade reagents supplied by Tokyo Kasel Kogyo Co., Ltd. Toluene and methylcyclohexane were used without further purification. (Trifluoromethyl)benzene and 1,2-dimethoxyethane were further purified by distillation in an Oldershaw distillation column with 30 plates. The physical properties of the materials used are listed in Table I.

## Results

The experimental total pressures $P$ for the three binary systems at 350 K are given in Table II and Figure 2. The experimental densities $\rho$ for the three binary mixtures at 298.15 K are shown in Table III. Figure 3 shows the molar excess liquid volumes $V^{E}$ calculated from

$$
\begin{equation*}
V^{E}=V-\left(x_{1} V_{1}+x_{2} V_{2}\right) \tag{1}
\end{equation*}
$$

and correlated with the following equation:

$$
\begin{equation*}
v^{E}=x_{1} x_{2}\left[\alpha+\beta\left(x_{1}-x_{2}\right)\right] \tag{2}
\end{equation*}
$$

